

# Synchronization of groups of coupled oscillators with sparse connections

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**Abstract** – Synchronization of groups of coupled oscillators with sparse connections are explored. It is found that different topologies of intergroup couplings may lead to different synchronizability. In the strong-coupling limit, an analytical treatment and criterion is proposed to judge the synchronization between communities of oscillators, and an optimal connection scheme for the group synchronization is given. By varying the intergroup and intragroup coupling strengths, different synchronous phases, *i.e.*, the unsynchronized state, intragroup synchronization, intergroup synchronization, and global synchronization are revealed. The present discussions and results can be applied to study the pattern formation and synchronization of coupled spatiotemporal systems.

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Synchronization, as a universal cooperative behavior and a fundamental mechanism of self-organizations in nature, has been extensively studied in recent years [1–6]. The synchronization of spatiotemporal systems have attracted much attention, and numerous interesting bifurcations and on-off intermittent behaviors are found [7]. The recent focus on complex networks has given much insight on the relation between the topological properties of coupled oscillators and collective dynamics [8]. The structural properties of a complex network may significantly affect the dynamical processes. It has been found that networks with a homogeneous distribution of connectivity are more synchronizable than heterogeneous ones (*e.g.*, scale-free networks) [9]. Synchronization can also be effectively improved on weighted networks [10].

Networks with community structures arise ubiquitously in social and biological sciences [11]. On such networks, elements in a society tend to form groups according to their social characteristics, and connections among elements in one group are dense while connections among elements in different groups are typically sparse. It is of potential interest to investigate the synchronization between two groups (we call it *group synchronization*, GrS)

when only a few intergroup connections are considered. Even for a given number of intergroup links, there can be many different ways of connections between two groups. The synchronization dynamics between two communities may be completely different for different network topologies. Therefore it is important to study this issue and give theoretical approaches or criteria for the emergence of GrS for a given coupled community network. This should be helpful in understanding the dynamical processes and their implications in relating to biological or social issues about community networks.

In this letter, we explore the synchronization of communities of coupled oscillators with sparse connections. It is found that different topologies of intergroup couplings may lead to different synchronizability. In the strong-coupling limit, we propose an analytical treatment and criterion to judge the synchronization between communities of oscillators, and an optimal connection scheme for the group synchronization is given. By varying the intergroup and intragroup coupling strengths, different synchronous phases, *i.e.*, the unsynchronized state, intragroup synchronization, intergroup synchronization, and global synchronization are revealed. Our exploration reveals the significant role of network topology in governing the global dynamics. The present discussions can be applied

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to study the pattern formation and synchronization of coupled spatiotemporal systems.

The dynamics on a given network can be written as

$$\dot{\vec{X}} = \vec{F}(\vec{X}) + \Gamma \otimes C \vec{X}, \quad (1)$$

where  $\vec{X} = (\vec{x}_1, \vec{x}_2, \dots, \vec{x}_N)$  denotes dynamical variables,  $\vec{F} = (\vec{f}_1, \vec{f}_2, \dots, \vec{f}_N)$  represents the nonlinear functions governing the dynamics of oscillators, and  $N$  is the number of nodes, and  $\Gamma: R^m \rightarrow R^m$  describes the coupling scheme among the nodes.  $C$  is the adjacent matrix representing the links among elements in a network, and  $\sum_{j=1}^N C_{ij} = 0$ , with the largest eigenvalue  $\lambda_1^C = 0$  ( $\lambda_1 \geq \lambda_2 \geq \dots \geq \lambda_N$ ).

For simplicity, we consider a network composed of two identical groups, with each group being  $N$  chaotic oscillators (labelled as  $i = 1, 2, \dots, N$ ) with nearest-neighbor coupling  $\varepsilon$ . Between two groups, there are  $n_l < N$  connections with the intergroup coupling strength  $r$ , which are labelled as  $L_1, L_2, \dots, L_{n_l}$ . A network with the this structure is marked as  $2 * N(L_1, L_2, \dots, L_{n_l})$ . We use indices  $i$  and  $i'$  to represent the corresponding nodes in two groups.

In the following discussions we apply the typical three-dimensional Rossler oscillator as our nodes:  $\dot{x} = -(y + z)$ ,  $\dot{y} = x + ay$ ,  $\dot{z} = b + z(x - c)$ , and each node is linearly coupled to others through the  $y$  component. Throughout this letter we set  $a = 0.15, b = 0.4, c = 8.5$ . The results we observed for the Rossler oscillators in this letter can be naturally generalized to other cases of node dynamics.

We first numerically study the dynamics for the case of  $N = 8$  and  $n_l = 2$ . Dynamical equations eqs. (1) are numerically integrated in terms of Runge-Kutta algorithm with the time step  $\Delta t = 0.001$ . The synchronization between two oscillators  $i$  and  $j$  can be measured by introducing the distance of two states as  $d_{ij} = \lim_{T \rightarrow \infty} (T^{-1} \int_0^T \|\vec{x}_i(t) - \vec{x}_j(t)\| dt)$ , where  $\|\cdot\|$  can be the Euclidean distance. GrS is achieved if  $d_{ii'} = 0$  but  $d_{ij} \neq 0$  for all  $i, j, i'$ . Intragroup synchronization occurs when  $d_{ij} = 0$  and  $d_{i'j'} = 0$  but  $d_{ii'} \neq 0$ . If for all sites in two groups satisfy  $d_{ij} = 0$ ,  $d_{i'j'} = 0$ , and  $d_{ii'} = 0$ , global synchronization can be achieved.

The synchronous dynamics for a network  $2 * 8(2, 7)$  in fig. 1(a) are presented in fig. 1(b). We keep the coupling ratio  $r/\varepsilon = 0.8$  and increase both  $r$  and  $\varepsilon$ . It can be found that GrS takes place prior to intragroup synchronization. For a different network topology  $2 * 8(2, 4)$  (see fig. 1(c)), the synchronization dynamics become completely different, as shown in fig. 1(d), where  $r/\varepsilon = 1$ . It can be found that GrS does not occur prior to complete synchronization, in fact they are reached at the same critical coupling strength. The contrasting results reveal that different network structures may lead to different synchronization phenomena, and the structural properties have a great impact on the synchronous dynamics. Therefore an important question naturally arises: what kind of networks can support the group synchronization? Furthermore, which structure can lead to an optimized GrS? It is thus very

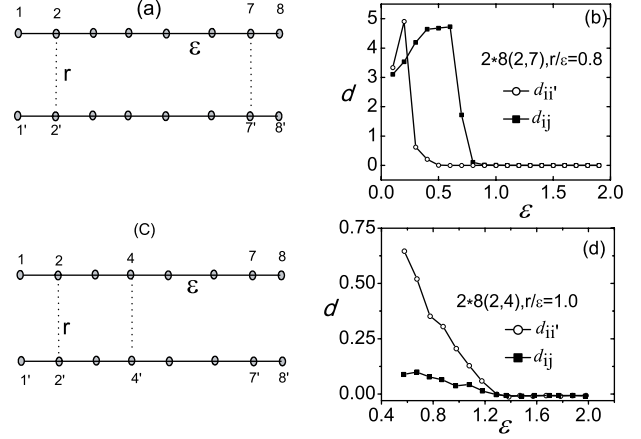


Fig. 1: (a) The topological graph of the network  $2 * 8(2, 7)$ . (b) The relation between the average distance and the coupling strength  $\varepsilon$  for  $r/\varepsilon = 0.8$ . (c) The topological graph of networks with two coupled chains  $2 * 8(2, 4)$ . (d) The relation between the average distance and the coupling strength on the network at  $r/\varepsilon = 1.0$ .

important to ravel the condition for the emergence of group synchronization.

The stability of complete synchronization (CS) can be well solved by applying the master stability function scheme [12]. CS is determined by the stability of the transversal subsystem. For two sparcely coupled communities, GrS can be considered as the CS between two groups, but it is difficult to give an eminent analytical expression because each group is a high-dimensional dynamical system. Recently, the partial synchronization (PaS) of coupled oscillators was explored theoretically and numerically [13], and analytical criterion for the emergence of PaS has been proposed. Obviously, the topology of two coupled groups satisfies the mirror symmetry, hence the treatment of PaS may be well applied to theoretical analysis of GrS.

The mirror symmetry of the network topology implies the invariance of the adjacent matrix  $C$  under the similarity transformation:  $F_{2N} C F_{2N}^{-1} = C$ , where  $F_{2N}$  is a counteridentity matrix satisfying  $F_{i, 2N+1-i} = 1$ ,  $i = 1, 2, \dots, 2N$  and other elements are all 0. Furthermore the adjacent matrix can be block diagonalized by constructing a similarity transformation  $S$ :

$$M = S^{-1} C S = \begin{pmatrix} A & 0 \\ 0 & B \end{pmatrix},$$

where

$$S = \begin{pmatrix} I_N & F_N \\ I_N & -F_N \end{pmatrix},$$

and  $I_N$  is the  $N$ -order identity matrix. By introducing new variables  $(\vec{G}_l, \vec{G}_r) = S(\vec{X}, \vec{X}')$ , where  $\vec{G}_l = \vec{X} + \vec{X}'$ ,  $\vec{G}_r = \vec{X} - \vec{X}'$ , and  $\vec{G}_l = (\vec{g}_l^1, \vec{g}_l^2, \dots, \vec{g}_l^N)$  and  $\vec{G}_r = (\vec{g}_r^1, \vec{g}_r^2, \dots, \vec{g}_r^N)$ . The similarity matrix  $S$  does not change the eigenvalues of the matrix  $C$ . It only redistributes them

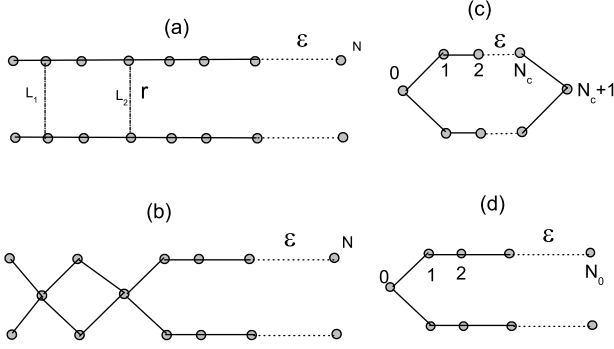


Fig. 2: The whole system is divided into  $N_c$ -close-block and  $N_o$ -open-block in the condition of infinite connection strength between groups.

to two blocks  $A$  and  $B$ . In fact,  $\vec{G}_l$  denotes the dynamics of the globally synchronous (Gls) manifold, and  $\vec{G}_r$  denotes the dynamics on the GrS manifold. Then the system can be divided into two parts. It is obvious that the GrS state is the solution of eq. (1). It is important to study the stability of the GrS state. The dynamics near the GrS state can be written as

$$\delta \dot{\vec{G}}_l = [D\vec{F}(\vec{G}_l) + \Gamma \otimes A] \delta \vec{G}_l, \quad (2)$$

$$\delta \dot{\vec{G}}_r = [D\vec{F}(\vec{G}_r) + \Gamma \otimes B] \delta \vec{G}_r, \quad (3)$$

where

$$[\Gamma \otimes A(\delta \vec{G}_l)]^i = \varepsilon \Gamma (\delta G_l^{i+1} + \delta G_l^{i-1} - 2\delta G_l^i). \quad (4)$$

For oscillators without intergroup couplings,

$$[\Gamma \otimes B(\delta \vec{G}_r)]^i = \varepsilon \Gamma (\delta G_r^{i+1} + \delta G_r^{i-1} - 2\delta G_r^i), \quad (5)$$

$i = 1, 2, \dots, N$ . For oscillators coupled with intergroup links:

$$[\Gamma \otimes B(\delta \vec{G}_r)]^j = \varepsilon \Gamma (\delta G_r^{j+1} + \delta G_r^{j-1}) - 2(\varepsilon + r) \Gamma \delta G_r^j \quad (6)$$

According to the criterion for the emergence of PaS [13], the criterion for the emergence of GrS is

$$\lambda_2^l > \lambda_1^r, \quad (7)$$

where  $\lambda_2^l$  is the second largest eigenvalue of  $A$ , and  $\lambda_1^r$  is the largest eigenvalue of  $B$ .

Let us further approach analytical results for the criterion of GrS. In the strong intergroup coupling limit, *i.e.*,  $r \rightarrow \infty$ , as shown in fig. 2(a), two directly coupled oscillators can be regarded as a common node of two groups (see fig. 2(b)). Therefore two coupled chains can be divided into several parts with two types of fundamental topologies: the ring block with  $2(N_c + 1)$  units (see fig. 2(c)) and the chain block with  $2N_o + 1$  oscillators (see fig. 2(d)). The possibility of GrS depends crucially on the synchronizability of these fundamental blocks. Note

that synchronizability for both fundamental structures can be theoretically analyzed, hence one may obtain analytical results by combining the criterion (7) for GrS.

In the strong-coupling limit, for the dynamics  $\vec{G}_l = \vec{X} + \vec{X}'$  on the GlS manifold, the eigenvalues of the connection matrix can be worked out as

$$\lambda_q^l = -4\varepsilon \sin^2(q/2), \quad q = m\pi/N, \quad (8)$$

where  $m = 0, 1, 2, \dots, N - 1$ .

For the dynamics  $\vec{G}_r = \vec{X} - \vec{X}'$  on the GrS manifold in the regime  $r \rightarrow \infty$ , the eigenvalues of the connection matrix for a ring block are given by

$$\lambda_q^r = -4\varepsilon \sin^2(q/2), \quad q = m\pi/(N_c + 1). \quad (9)$$

Here  $m = 1, 2, \dots, N_c$ .

For the case of chain blocks, the eigenvalues of the connection matrix can also be analytically worked out as

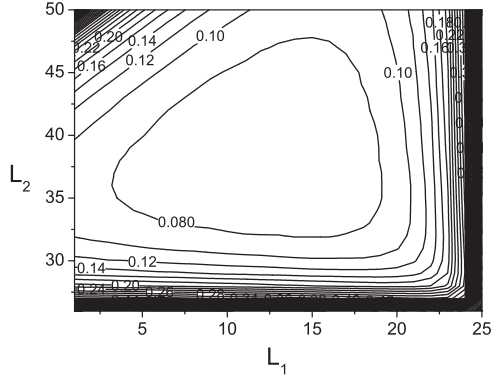
$$\lambda_q^r = -4\varepsilon \sin^2(q/2), \quad q = (2m - 1)\pi/(2N_o + 1), \quad (10)$$

where  $m = 1, 2, \dots, N_o$ .

The second largest eigenvalue for the GlS manifold is  $\lambda_2^l = -4\varepsilon \sin^2(\pi/2N)$ . For the GrS manifold, the largest eigenvalue for the transversal manifold is  $\lambda_1^r = -4\varepsilon \sin^2(\pi/2(N_c + 1))$  for the ring block. For the chain structure, one has  $\lambda_1^r = -4\varepsilon \sin^2(\pi/(2N_o + 1))$ . By applying the criterion (7) for GrS, one can naturally give rise to the following conclusions:

- For both types of blocks, the larger the size of a block is, the larger  $\lambda_1^r$  is, and the more difficult for  $\lambda_2^l > \lambda_1^r$ . Therefore GrS becomes difficult if the size of one of the blocks in the network becomes very large.
- If there exists a chain block with  $N_o \geq N/2$ , *i.e.*, its size exceeds one half of the network size, one has  $\lambda_2^l < \lambda_1^r$ , *i.e.*, GrS will not be observed.
- A chain block with the size  $n$  has the same eigenvalue  $\lambda_1^r$  and the same synchronization ability as a ring block with size  $2n$ .
- The minimum number of intergroup links is  $n_l = 2$ , *i.e.*, GrS cannot be achieved when there is only one intergroup link. For the case of  $n_l = 2$ , the optimal way for GrS is the topology that includes form two chains with the size  $n = (N - 2)/4$  and one loop block with the size  $(N - 2)/2$ .

In terms of the above theoretical arguments, it is not difficult to understand the synchronization phenomena in the aforementioned networks  $2 * 8(2, 7)$  and  $2 * 8(2, 4)$ . For the network  $2 * 8(2, 7)$ , it is possible to achieve GrS when  $r/\varepsilon > 0.42$ . While for the structure  $2 * 8(2, 4)$ , there is one chain with  $N_o = 4 = N/2$ , and  $\lambda_2^l = [-4\varepsilon \sin^2(\pi/16)] < \lambda_1^r = [-4\varepsilon \sin^2(\pi/18)]$ . So GrS cannot be reached. Note



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